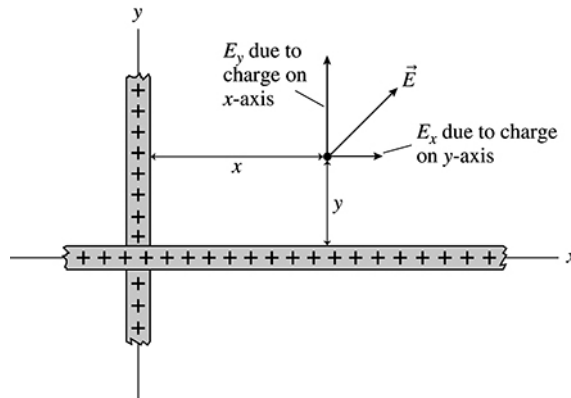


**27.38. Model:** The fields are those of two infinite lines of charge with linear charge density  $\lambda$ .  
**Visualize:**



**Solve:** The field at distance  $r$  from an infinite line of charge is

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{2\lambda}{r} \hat{r}$$

It points straight away from the line. With two perpendicular lines, the field due to the line along the  $x$ -axis points in the  $y$ -direction and depends inversely on distance  $y$ . Similarly, the field due to the line along the  $y$ -axis points in the  $x$ -direction and depends inversely on distance  $x$ . That is

$$E_x = \frac{1}{4\pi\epsilon_0} \frac{2\lambda}{x} \quad E_y = \frac{1}{4\pi\epsilon_0} \frac{2\lambda}{y}$$

$$\Rightarrow \vec{E} = E_x \hat{i} + E_y \hat{j} = \frac{2\lambda}{4\pi\epsilon_0} \left( \frac{1}{x} \hat{i} + \frac{1}{y} \hat{j} \right)$$

The field strength at this point in space is

$$E = \sqrt{E_x^2 + E_y^2} = \frac{2\lambda}{4\pi\epsilon_0} \sqrt{1/x^2 + 1/y^2}$$

**27.45. Model:** Assume that the ring of charge is thin and that the charge lies along circle of radius  $R$ .  
**Solve:** (a) From Example 27.5, the on-axis field of a ring of charge  $Q$  and radius  $R$  is

$$E_z = \frac{1}{4\pi\epsilon_0} \frac{zQ}{(z^2 + R^2)^{3/2}}$$

For the field to be maximum at a particular value of  $z$ ,  $dE/dz = 0$ . Taking the derivative,

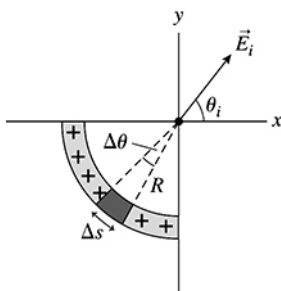
$$\begin{aligned} \frac{dE}{dz} &= \frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{(z^2 + R^2)^{3/2}} - \frac{z(3/2)(2z)}{(z^2 + R^2)^{5/2}} \right] = 0 \Rightarrow \frac{1}{(z^2 + R^2)^{3/2}} = \frac{3z^2}{(z^2 + R^2)^{5/2}} \\ &\Rightarrow 1 = \frac{3z^2}{z^2 + R^2} \Rightarrow z = \frac{R}{\sqrt{2}} \end{aligned}$$

(b) The field strength at the point  $z = R/\sqrt{2}$  is

$$(E_z)_{\max} = \frac{Q}{4\pi\epsilon_0} \frac{(R/\sqrt{2})}{\left[ (R/\sqrt{2})^2 + R^2 \right]^{3/2}} = \frac{2}{3\sqrt{3}} \frac{Q}{4\pi\epsilon_0 R^2}$$

**27.47. Model:** Assume that the quarter-circle plastic rod is thin and that the charge lies along the quarter-circle of radius  $R$ .

**Visualize:**



The origin of the coordinate system is at the center of the circle. Divide the rod into many small segments of charge  $\Delta q$  and arc length  $\Delta s$ .

**Solve:** (a) Segment  $i$  creates a small electric field  $\vec{E}_i$  at the origin with two components:

$$(E_i)_x = E_i \cos \theta_i \quad (E_i)_y = E_i \sin \theta_i$$

Note that the angle  $\theta_i$  depends on the location of the segment  $i$ . Now all segments are at distance  $r_i = R$  from the origin, so

$$E_i = \frac{1}{4\pi\epsilon_0} \frac{\Delta q}{r_i^2} = \frac{1}{4\pi\epsilon_0} \frac{\Delta q}{R^2}$$

The linear charge density of the rod is  $\lambda = Q/L$ , where  $L$  is the rod's length ( $L = \text{quarter-circumference} = \pi R/2$ ). This allows us to relate charge  $\Delta q$  to the arc length  $\Delta s$  through

$$\Delta q = \lambda \Delta s = \left(\frac{Q}{L}\right) \Delta s = \left(\frac{2Q}{\pi R}\right) \Delta s$$

Using  $\Delta s = R\Delta\theta$ , the components of the electric field at the origin are

$$(E_i)_x = \frac{1}{4\pi\epsilon_0} \frac{\Delta q}{R^2} \cos \theta_i = \frac{1}{4\pi\epsilon_0} \frac{1}{R^2} \left(\frac{2Q}{\pi R}\right) R \Delta\theta \cos \theta_i = \frac{1}{4\pi\epsilon_0} \left(\frac{2Q}{\pi R^2}\right) \Delta\theta \cos \theta_i$$

$$(E_i)_y = \frac{1}{4\pi\epsilon_0} \frac{\Delta q}{R^2} \sin \theta_i = \frac{1}{4\pi\epsilon_0} \frac{1}{R^2} \left(\frac{2Q}{\pi R}\right) R \Delta\theta \sin \theta_i = \frac{1}{4\pi\epsilon_0} \left(\frac{2Q}{\pi R^2}\right) \Delta\theta \sin \theta_i$$

(b) The  $x$ - and  $y$ -components of the electric field for the entire rod are the integrals of the expressions in part (a) from

$\theta = 0$  rad to  $\theta = \pi/2$ . We have

$$E_x = \frac{1}{4\pi\epsilon_0} \left(\frac{2Q}{\pi R^2}\right) \int_0^{\pi/2} \cos \theta d\theta \quad E_y = \frac{1}{4\pi\epsilon_0} \left(\frac{2Q}{\pi R^2}\right) \int_0^{\pi/2} \sin \theta d\theta$$

(c) The integrals are

$$\int_0^{\pi/2} \sin \theta d\theta = [-\cos \theta]_0^{\pi/2} = -\left(\cos \frac{\pi}{2} - \cos 0\right) = +1 \quad \int_0^{\pi/2} \cos \theta d\theta = [\sin \theta]_0^{\pi/2} = \sin \frac{\pi}{2} - \sin 0 = +1$$

The electric field is

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{2Q}{\pi R^2} (\hat{i} + \hat{j})$$

**27.53. Model:** Assume that the electric field inside the capacitor is constant, so constant-acceleration kinematic equations apply.

**Visualize:** Please refer to Figure P27.53.

**Solve:** (a) The force on the electron inside the capacitor is

$$\vec{F} = m\vec{a} = q\vec{E} \Rightarrow \vec{a} = \frac{q\vec{E}}{m}$$

Because  $\vec{E}$  is directed upward (from the positive plate to the negative plate) and  $q = -1.60 \times 10^{-19} \text{ C}$ , the acceleration of the electron is downward. We can therefore write the above equation as simply  $a_y = qE/m$ . To determine  $E$ , we must first find  $a_y$ . From kinematics,

$$\begin{aligned} x_1 &= x_0 + v_{0x}(t_1 - t_0) + \frac{1}{2}a_x(t_1 - t_0)^2 \Rightarrow 0.040 \text{ m} = 0 \text{ m} + v_0 \cos 45^\circ(t_1 - t_0) + 0 \text{ m} \\ \Rightarrow (t_1 - t_0) &= \frac{(0.040 \text{ m})}{(5.0 \times 10^6 \text{ m/s})\cos 45^\circ} = 1.1314 \times 10^{-8} \text{ s} \end{aligned}$$

Using the kinematic equations for the motion in the  $y$ -direction,

$$\begin{aligned} v_{1y} &= v_{0y} + a_y \left( \frac{t_1 - t_0}{2} \right) \Rightarrow 0 \text{ m/s} = v_0 \sin 45^\circ + \left( \frac{qE}{m} \right) \left( \frac{t_1 - t_0}{2} \right) \\ \Rightarrow E &= -\frac{2 m v_0 \sin 45^\circ}{q(t_1 - t_0)} = -\frac{2(9.1 \times 10^{-31} \text{ kg})(5.0 \times 10^6 \text{ m/s})\sin 45^\circ}{(-1.60 \times 10^{-19} \text{ C})(1.1314 \times 10^{-8} \text{ s})} = 3550 \text{ N/C} = 3.6 \times 10^3 \text{ N/C} \end{aligned}$$

(b) To determine the separation between the two plates, we note that  $y_0 = 0 \text{ m}$  and  $v_{0y} = (5.0 \times 10^6 \text{ m/s})\sin 45^\circ$ , but at  $y = y_1$ , the electron's highest point,  $v_{1y} = 0 \text{ m/s}$ . From kinematics,

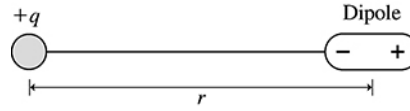
$$\begin{aligned} v_{1y}^2 &= v_{0y}^2 + 2a_y(y_1 - y_0) \Rightarrow 0 \text{ m}^2/\text{s}^2 = v_0^2 \sin^2 45^\circ + 2a_y(y_1 - y_0) \\ \Rightarrow (y_1 - y_0) &= -\frac{v_0^2 \sin^2 45^\circ}{2a_y} = -\frac{v_0^2}{4a_y} \end{aligned}$$

From part (a),

$$\begin{aligned} a_y &= \frac{qE}{m} = \frac{(-1.60 \times 10^{-19} \text{ C})(3550 \text{ N/C})}{9.1 \times 10^{-31} \text{ kg}} = -6.242 \times 10^{14} \text{ m/s}^2 \\ \Rightarrow y_1 - y_0 &= -\frac{(5.0 \times 10^6 \text{ m/s})^2}{4(-6.242 \times 10^{14} \text{ m/s}^2)} = 0.010 \text{ m} = 1.0 \text{ cm} \end{aligned}$$

This is the height of the electron's trajectory, so the minimum spacing is 1.0 cm.

**27.60. Model:** The electric field at the dipole's location is that of the ion with charge  $q$ .  
**Visualize:**



**Solve:** (a) We have  $p = \alpha E$ . The units of  $\alpha$  are the units of  $p/E$  and are

$$\frac{\text{C m}}{\text{N/C}} = \frac{\text{C}^2 \text{ m}}{\text{N}} = \frac{\text{C}^2 \text{ s}^2}{\text{kg}}$$

(b) The electric field due to the ion at the location of the dipole is

$$E_{\text{at dipole}} = \left( \frac{1}{4\pi\epsilon_0} \frac{q}{r^2}, \text{ away from } q \right) = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{i}$$

Because  $\vec{p} = \alpha \vec{E}$ , the induced dipole moment is

$$\vec{p} = \alpha \left( \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \right) \hat{i}$$

From Equation 27.11, the electric field produced by the dipole at the location of the ion is

$$\vec{E}_{\text{dipole}} = \frac{1}{4\pi\epsilon_0} \frac{2\vec{p}}{r^3} = \frac{1}{4\pi\epsilon_0} \left( \frac{2}{r^3} \right) \alpha \left( \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \right) \hat{i} = \left( \frac{1}{4\pi\epsilon_0} \right)^2 \left( \frac{2q\alpha}{r^5} \right) \hat{i}$$

The force the dipole exerts on the ion is

$$\vec{F}_{\text{dipole on ion}} = q\vec{E}_{\text{dipole}} = \left( \frac{1}{4\pi\epsilon_0} \right)^2 \left( \frac{2q^2\alpha}{r^5} \right) \hat{i}$$

According to Newton's third law,  $\vec{F}_{\text{dipole on ion}} = -\vec{F}_{\text{ion on dipole}}$ . Therefore,

$$\vec{F}_{\text{ion on dipole}} = \left( \left( \frac{1}{4\pi\epsilon_0} \right)^2 \frac{2q^2\alpha}{r^5}, \text{ toward ion} \right)$$